

A hyperelastic model for cardiac deformation applied to multimodal registration of MRI and optical mapping data

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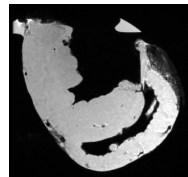
Table of contents

- 1 Introduction
- 2 Hyperelastic model
- 3 Numerical simulation
- 4 Conclusion

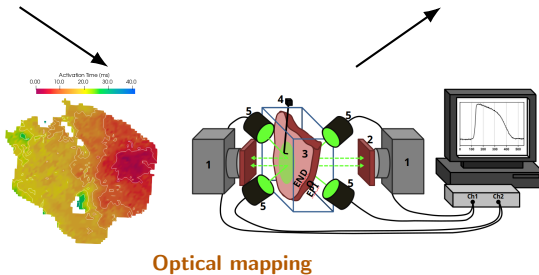
Motivation



MRI



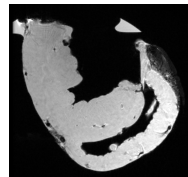
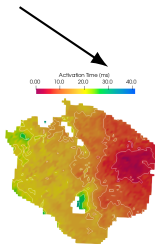
High-resolution MRI (9.4T)



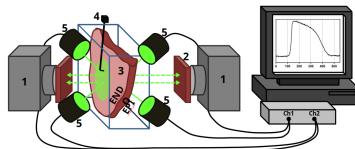
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MRI
Geometry



High-resolution MRI (9.4T)
Microstructure - Fibers

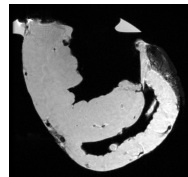
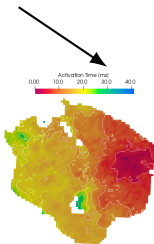


Optical mapping
Electrical activation map

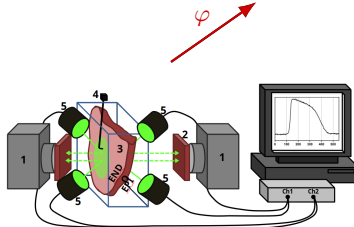
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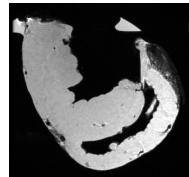
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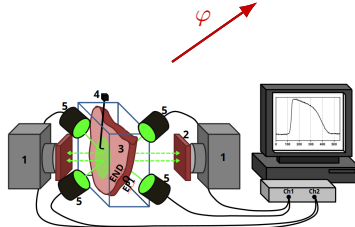
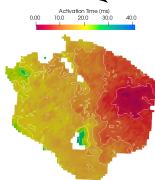


MRI
Geometry

Goal: combine **structural** and **functional** information to better understand **arrhythmia** , responsible for a lot of cardiovascular deaths.



High-resolution MRI (9.4T)
Microstructure - Fibers



Optical mapping
Electrical activation map

Key idea

We want to deform the MRI geometry onto the optical mapping geometry.

$$\varphi = \text{id} + \mathbf{u} : \Omega_0 \rightarrow \Omega$$

- Ω_0 : reference domain (resting position);
- Ω : deformed domain (stretched position);
- \mathbf{u} : displacement field.

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- Ω_0 : reference domain (resting position);
- Ω : deformed domain (stretched position);
- \mathbf{u} : displacement field.

The displacement \mathbf{u} is obtained by solving a nonlinear hyperelastic equilibrium problem:

$$\begin{cases} -\text{div}(\mathbf{P}(\mathbf{u})) &= \mathbf{f} & \text{in } \Omega_0, \\ \mathbf{P}(\mathbf{u}) \cdot \mathbf{n} &= \mathbf{g} & \text{on } \Gamma_N, \\ \mathbf{u} &= \mathbf{0} & \text{on } \Gamma_D, \end{cases}$$

where

- $\partial\Omega_0 = \Gamma_D \cup \Gamma_N$, Γ_D is the Dirichlet boundary and Γ_N is the Neumann boundary
- \mathbf{P} is the first Piola–Kirchhoff stress tensor,
- \mathbf{n} the outward unit normal vector to $\partial\Omega_0$,
- \mathbf{f} denotes the body force in Ω_0 and \mathbf{g} denotes the prescribed traction on Γ_N .

Table of contents

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Kinematics and hyperelastic definitions

The **deformation** is described by the map

$$\varphi : \Omega_0 \rightarrow \mathbb{R}^d, \quad \varphi(\mathbf{X}) = \mathbf{X} + \mathbf{u}(\mathbf{X}), \quad \mathbf{X} \in \Omega_0,$$

where $\mathbf{u} : \Omega_0 \rightarrow \mathbb{R}^d$ is the **displacement field**.

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The main quantities are

- $\mathbf{F} := \nabla \varphi = \text{Id} + \nabla \mathbf{u}$ is the **deformation gradient**,
- $J := \det \mathbf{F}$ measures the **local change of volume**,
- $\mathbf{C} := \mathbf{F}^T \mathbf{F}$ is the **right Cauchy–Green tensor** and measures strain in Ω_0 .

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Hyperelastic material

A hyperelastic material is fully described by a **stored energy density** W .

The associated **stress tensor** is defined by

$$\mathbf{P} = \frac{\partial W}{\partial \mathbf{F}}.$$

This framework is well suited to **large deformations** and **soft biological tissues**.

Examples of hyperelastic materials

- **Saint Venant–Kirchhoff** → a simple extension of linear elasticity
- **Neo-Hookean** → simple isotropic model for rubber-like soft materials
- **Mooney–Rivlin** → an enriched Neo-Hookean model

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In this work

We use a nearly-incompressible **Holzappel-Gasser-Ogden type energy** to model cardiac tissue, with **reduced invariants**.

Invariants

The HGO model is written with respect to scalar invariants of the tensor \mathbf{C} .

Principal invariants

$$\det(\mathbf{C} - \lambda \text{Id}) = \sum_{p=0}^d (-1)^p I_p \lambda^{d-p}, \quad I_0 := 1.$$

In particular, in 3D,

- $I_1 := \text{tr}(\mathbf{C})$ measures the amount of stretch;
- $I_2 := \frac{1}{2}(I_1^2 - \text{tr}(\mathbf{C}^2))$ measures area changes;
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Fiber invariant

To represent the local fiber direction, we denote:

- $\mathbf{a} : \Omega_0 \rightarrow \mathbb{R}^d$ the tangential unit vector in the fiber direction;
- the additional invariant that measures the squared stretch in the fiber direction,

$$I_4 := \mathbf{a}^T \mathbf{C} \mathbf{a} = \|\mathbf{F} \mathbf{a}\|^2.$$

The Holzapfel–Gasser–Ogden energy for fiber-reinforced cardiac tissue

Assumptions on the material

- **Nearly incompressible:** local volume changes are strongly penalized, so $J \simeq 1$;
- **Transverse isotropy** induced by the fiber field $\mathbf{a} : \Omega_0 \rightarrow \mathbb{R}^d$.
- **Isochoric split:** shape changes and volume changes are modeled separately.

The hyperleastic energy is given by,

$$\begin{aligned}
 W(\mathbf{C}(\mathbf{u})) &= \underbrace{\widetilde{W}(\widetilde{\mathbf{I}}_1, \widetilde{\mathbf{I}}_2, \widetilde{\mathbf{I}}_4)}_{\text{isochoric part}} + \underbrace{W_{\text{vol}}(J)}_{\text{volumetric part}} \\
 &= \underbrace{W_{\text{lin}}(\widetilde{\mathbf{I}}_1, \widetilde{\mathbf{I}}_2) + W_{\text{iso}}(\widetilde{\mathbf{I}}_1) + W_{\text{fib}}(\widetilde{\mathbf{I}}_4)} + W_{\text{vol}}(J).
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Since near-incompressibility is handled separately by $W_{\text{vol}}(J)$, the isochoric part is written in terms of reduced invariants given by

$$\widetilde{I}_1 = J^{-2/d} I_1, \quad \widetilde{I}_2 = J^{-4/d} I_2, \quad \widetilde{I}_4 = J^{-2/d} I_4.$$

Interpretation of the energy terms and material parameters

$$W(\mathbf{C}(\mathbf{u})) = W_{\text{lin}}(\tilde{I}_1, \tilde{I}_2) + W_{\text{iso}}(\tilde{I}_1) + W_{\text{fib}}(\tilde{I}_4) + W_{\text{vol}}(J)$$

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 - ▶ C_0 : amplitude of the isotropic exponential contribution
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- A volume penalizing term: $W_{vol}(J) = \kappa(J - 1) - \kappa \ln J$,
 - ▶ this term enforces near-incompressibility
 - ▶ κ : bulk modulus parameter

From stored energy to the equilibrium equations

We seek a displacement $\mathbf{u} : \Omega_0 \rightarrow \mathbb{R}^d$ such that

$$\mathbf{u} \in \arg \min_{\mathbf{v} \in V} \mathbf{E}(\mathbf{v}), \quad V := \left\{ \mathbf{v} \in H^1(\Omega_0; \mathbb{R}^d) : \mathbf{v} = \mathbf{0} \text{ on } \Gamma_D \right\}.$$

where the total energy is

$$\mathbf{E}(\mathbf{u}) = \underbrace{\int_{\Omega_0} W(\mathbf{C}(\mathbf{u})) \, dx}_{\text{hyperelastic energy}} - \underbrace{\int_{\Omega_0} \mathbf{f} \cdot \mathbf{u} \, dx - \int_{\Gamma_N} \mathbf{g} \cdot \mathbf{u} \, ds}_{\text{external forces}}.$$

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Instead, we seek \mathbf{u} such that

$$\mathbf{DE}(\mathbf{u}; \mathbf{v}) = 0 \quad \forall \mathbf{v} \in V. \quad (1)$$

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with $\mathbf{P}(\mathbf{u})$ denotes the **first Piola–Kirchhoff stress tensor**.

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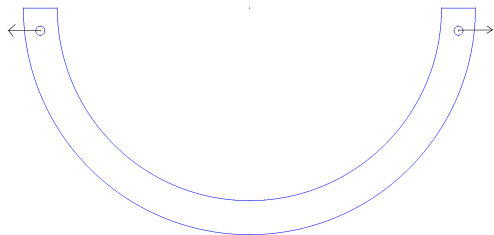
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 After integration by parts (1), we recover the full hyperelastic equilibrium system.

Table of contents

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Test case in 2D



	Symbol	Value
Inner radius	r_i	0.85
Outer radius	r_o	1.00
Holes' radius	ε	0.02

- Opposite horizontal tractions on the two holes:

$$\mathbf{g}_{left} = -g_0 \mathbf{e}_x, \quad \mathbf{g}_{right} = +g_0 \mathbf{e}_x.$$

- No body force is considered, $\mathbf{f} = 0$.
- A Dirichlet condition is imposed on 2% of the inner arc to remove rigid-body motions.

Finite element discretization using FreeFem++

The domain Ω is discretized with a **triangular mesh** in 2D.

- Maillage global : 10614 triangles

The displacement \mathbf{u} is approximated in a \mathbb{P}_1 **Lagrange finite element space**:

$$\mathbf{u}_h \in V_h.$$

Discrete equilibrium: Find $\mathbf{u}_h \in V_h$ such that

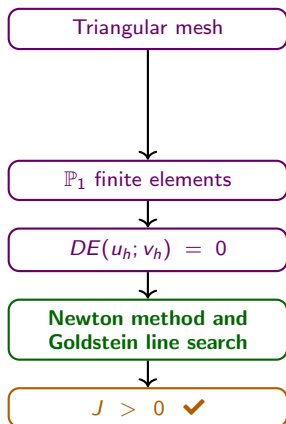
$$DE(\mathbf{u}_h; \mathbf{v}_h) = 0 \quad \forall \mathbf{v}_h \in V_h.$$

Admissibility constraint: We only consider deformations such that

$$J = \det(\nabla\varphi) > 0$$

to avoid element inversion.

Parameters	ν_1	C_0	C_1	C_2	C_3	κ
Values	1	10	0.5	15	0.9	800



Finite element discretization using FreeFem++

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- Maillage global : 10614 triangles
- Maillage local : 1782 triangles

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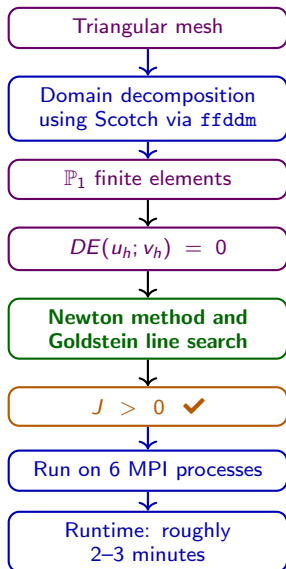
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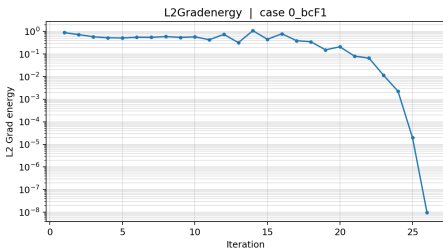
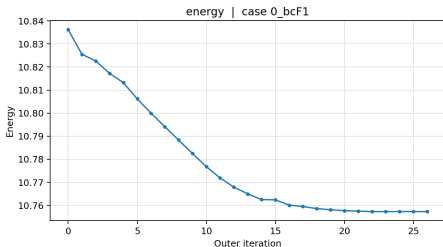
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Energy convergence indicators



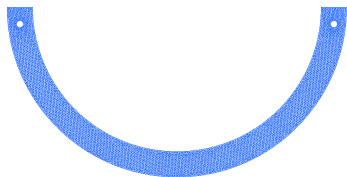
Results

- The energy decreases progressively during the iterations.
- The L^2 -norm of the energy gradient tends to zero at convergence.

Quasi-incompressible criteria :

- $J = \det(\nabla\varphi) > 0$ ✓
- $0.985 < J < 1.02$ ✓

Deformed domain



Reference domain

$$\varphi = \text{id} + \mathbf{u}$$



Deformed domain

Post-processing:

- The mesh is progressively deformed from the **reference** to the **deformed** configuration.
- No mesh element inversion occurs ✓
- **Observed volume variation:** 0.045%
- The **fiber field** is transported from the **reference** configuration to the **deformed** domain.

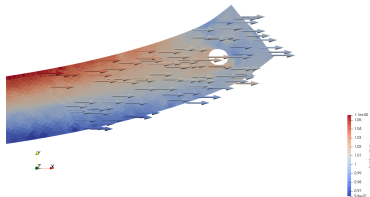


Table of contents

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Conclusion and next steps

What was established

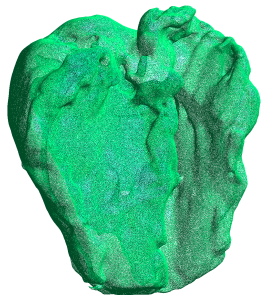
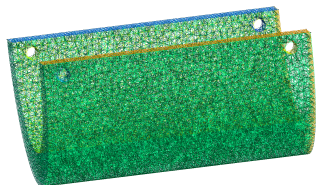
- Modeling the cardiac mechanical behavior using a Holzapfel–Ogden energy.
- Solving the 2D toy model validating the established code.

On going work

- Solve the direct problem on a 3D toy domains.
- Solving the associated inverse problem in 2D.

Next steps

- Solving the associated inverse problem in 3D.
- Move to the 3D case on segmentation of a heart tissue sample.



Thank you for your attention

Research interests

- **Numerical analysis:** high-order boundary conditions, error estimation, spectral problems, linear elasticity models, and nonlinear hyperelasticity with large deformations.
- **Scientific computing:** surface/volume finite elements, curved meshes, shape optimization, and registration methods.



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