

Asymptotic modeling of a random rough thin layer for electromagnetic wave scattering

CANUM 2026

02 juin 2026

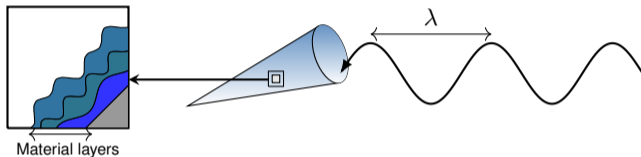
Pierre Boulogne^{1 2} Sonia Fliss² Laure Giovangigli² Justine Labat¹

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²POEMS, CNRS, INRIA, ENSTA, INSTITUT POLYTECHNIQUE DE PARIS, PALAISEAU, FRANCE

Context and motivation

- Context: Stacks of thin layers are used to improve electromagnetic stealth of objects they cover. **Small surface defects** may affect this stealth, hence their impact must be studied.



- Challenge: High computational costs due to the strongly **multi-scale** problem
- Proposed solution: Derivation of **effective conditions** replacing the complete stack which take roughness into account

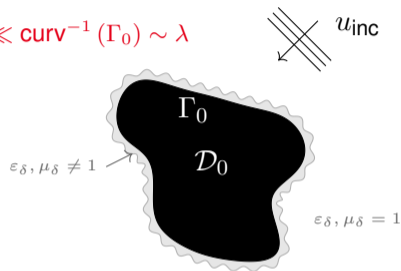
Thin layers: Bendali and Lemrabet (1996), Duruflé, Haddar, and Joly (2006)

Periodic setting: Abboud and Ammari (1996), Ammari and He (1998), Madureira and Valentin (2007), Delourme and Claeys (2012), Chamaillard (2016)

Random setting: Gérard-Varet (2008), Boucart (2023)

Model problem

$$\delta \ll \text{curv}^{-1}(\Gamma_0) \sim \lambda$$



- The total field u_δ satisfies the **2D Helmholtz equation**

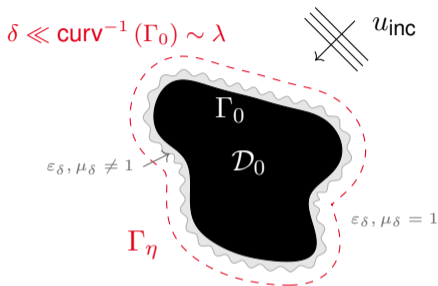
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Local parametrization

$$\varepsilon_\delta(s, \nu) = \varepsilon \left(\frac{s}{\delta}, \frac{\nu}{\delta} \right)$$

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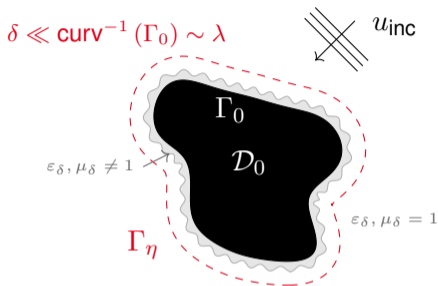
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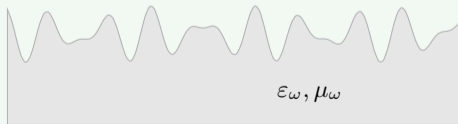
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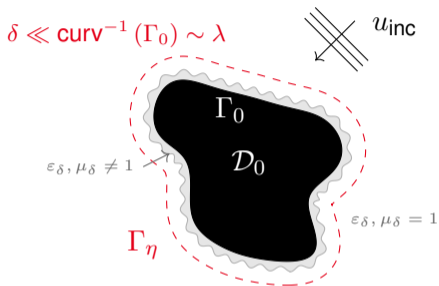
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Random Case

$(X, \omega) \mapsto (\varepsilon, \mu)_\omega(X, \cdot)$ is **stationary ergodic**



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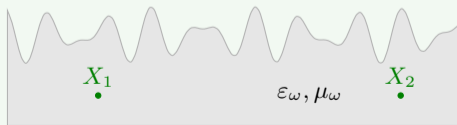
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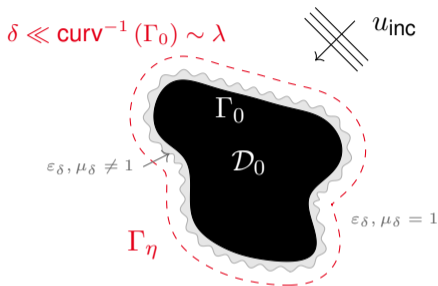
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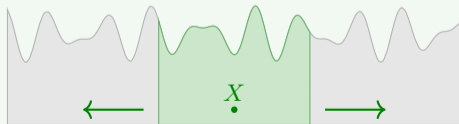
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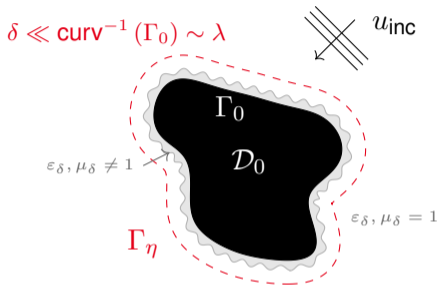
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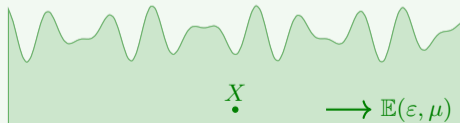
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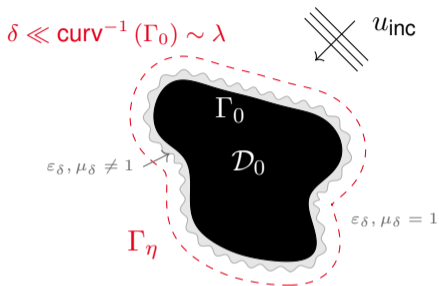
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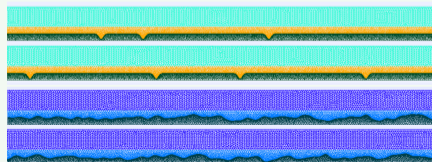
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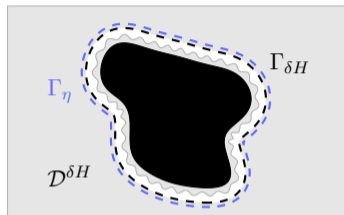
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Multi-scale expansions

- Near-field ansatz, $0 < \nu < \delta H$

$$u_{\omega, \delta}(s, \nu) = \sum_{n=0}^{\infty} \delta^n u_{n, \omega}^{\text{NF}} \left(s; \frac{s}{\delta}, \frac{\nu}{\delta} \right)$$



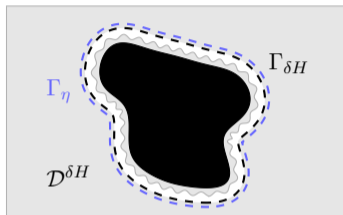
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- Far-field ansatz, $\delta H < \nu < \eta$

$$u_{\omega, \delta}(s, \nu) = \sum_{n=0}^{\infty} \delta^n \left[u_{n, \omega}^{\text{NF}} \left(s; \frac{s}{\delta}, \frac{\nu}{\delta} \right) + u_{n, \omega}^{\text{FF}}(s, \nu) \right]$$



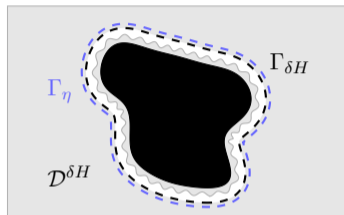
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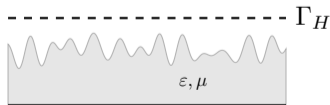
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$\lim_{Y \rightarrow +\infty} u_{n, \omega}^{\text{NF}}(\cdot; \cdot, Y) = 0$ and $(X, \omega) \mapsto u_{n, \omega}^{\text{NF}}(\cdot; X, \cdot)$ is stationary ergodic



Far-field problems

- The far field terms satisfy the Helmholtz equation

$$\text{(FFn)} \begin{cases} -\Delta u_{n,\omega}^{\text{FF}} - \kappa_0^2 u_{n,\omega}^{\text{FF}} = 0 & \text{in } \mathcal{D}^{\delta H} \\ u_{n,\omega}^{\text{FF}} - \delta_{n,0} u_{\text{inc}} & \text{is outgoing} \end{cases}$$

A condition on $\Gamma_{\delta H}$ is missing

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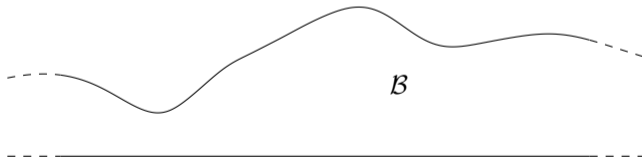
Condition on $\Gamma_{\delta H}$

This condition will be obtained by ensuring the existence of near-field terms satisfying our assumptions

Near-field problems

- The near fields satisfy the problems parametrized by s :

$$(\text{NFn}) \left\{ \begin{array}{l} -\nabla_{X,Y} \left(\frac{1}{\mu_\omega} \nabla_{X,Y} u_{n,\omega}^{\text{NF}}(s; \cdot, \cdot) \right) = \mathcal{F}_s \left[\left(u_{i,\omega}^{\text{NF}}(s; \cdot, \cdot) \right)_{i \leq n-1} \right] \quad \text{in } \mathcal{B} := \mathbb{R} \times \mathbb{R}_+ \\ [u_{n,\omega}^{\text{NF}}]_{\Gamma_H}(s) = -u_{n,\omega}^{\text{FF}}|_{\Gamma_{\delta H}}(s) \quad [\partial_Y u_{n,\omega}^{\text{NF}}]_{\Gamma_H}(s) = -\partial_\nu u_{n-1,\omega}^{\text{FF}}|_{\Gamma_{\delta H}}(s) \\ u_{n,\omega}^{\text{NF}}(s; \cdot, \cdot) = 0 \quad \text{on } \Sigma \\ (\omega, X) \mapsto u_{n,\omega}^{\text{NF}}(s; X, \cdot) \text{ is stationary} \end{array} \right.$$



Near-field problems - Random case

- The general corrector problem is:

(Boucart (2023))

$$(\mathbf{NF}) \begin{cases} -\operatorname{div}\left(\frac{1}{\mu_\omega} \nabla U_\omega\right) = F_\omega & \text{in } \mathcal{B} \\ [U_\omega]_{\Gamma_H} = \alpha_\omega^D & [\partial_Y U_\omega]_{\Gamma_H} = \alpha_\omega^N \\ U_\omega = 0 & \text{on } \Sigma \\ (\omega, X) \mapsto U_\omega(X, \cdot) & \text{is stationary} \end{cases}$$



The Dirichlet condition allows the use of a Hardy inequality in

$$V_0 := \left\{ \text{a.s. } v_\omega \in H_{\text{loc}}^1(\overline{\mathcal{B}}), v_\omega|_\Sigma = 0, (\omega, X) \mapsto v_\omega(X, \cdot) \text{ is stationary} \right. \\ \left. \text{and } \mathbb{E} \left[\int_0^{+\infty} |\nabla v|^2 + \frac{|v|^2}{1+Y^2} \right] < +\infty \right\}$$

in the form

$$\mathbb{E} \left[\int_0^\infty \frac{|v|^2}{1+Y^2} \right] \lesssim \mathbb{E} \left[\int_0^\infty |\nabla v|^2 \right]$$

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Provided that $\sqrt{1 + Y^2} F_\omega \in L^2(\Omega, L^2(\mathcal{B}))$, there exists a unique solution in $V_0 + \alpha_\omega^D \chi_{Y>H}$.

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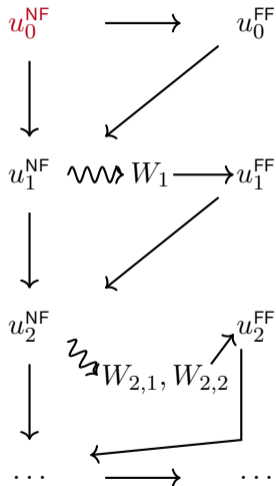


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When $F_\omega = 0$, $\lim_{Y \rightarrow \infty} U_\omega = 0$ iff a relation between α_ω^D and α_ω^N is satisfied.
In this case, we can simply show that:

$$\mathbb{E} \left[|U_\omega|^2 + Y^2 |\nabla U_\omega|^2 \right] \xrightarrow{Y \rightarrow \infty} 0$$

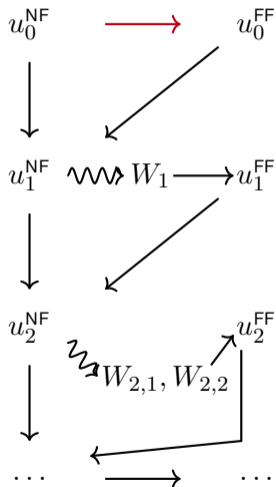
Formal recursive construction



- Problem satisfied by $u_{0,\omega}^{\text{NF}}$:

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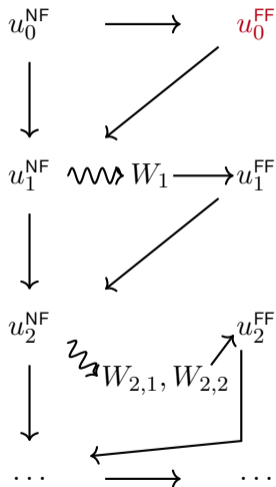
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Solution of (NF0)

$$\begin{aligned} u_{0,\omega}^{\text{NF}} + u_{0,\omega}^{\text{FF}}|_{\Gamma_{\delta H}} \chi_{Y>H} &= 0 \\ Y \rightarrow \infty \implies u_{0,\omega}^{\text{FF}} &= 0 \quad \text{on } \Gamma_{\delta H} \end{aligned}$$

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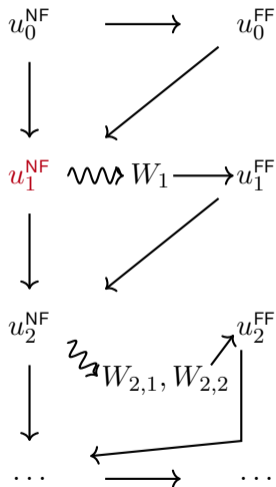
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$\implies u_{0,\omega}^{\text{NF}} = 0$ and $u_{0,\omega}^{\text{FF}}$ satisfies:

$$\text{Deterministic} \leftarrow \begin{cases} -\Delta u_{0,\omega}^{\text{FF}} - \kappa_0^2 u_{0,\omega}^{\text{FF}} = 0 & \text{in } \mathcal{D}^{\delta H} \\ u_{0,\omega}^{\text{FF}} = 0 & \text{on } \Gamma_{\delta H} \\ u_{0,\omega}^{\text{FF}} - u_{\text{inc}} & \text{is outgoing} \end{cases}$$

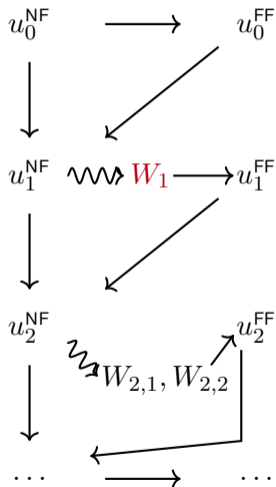
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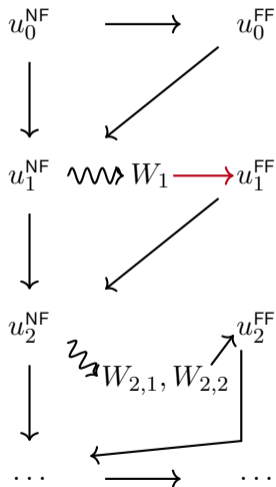
Formal recursive construction



- We can write $(u_{1,\omega}^{\text{NF}} + u_{1,\omega}^{\text{FF}} |_{\Gamma_{\delta H}} \chi_{Y>H}) = \partial_\nu u_{0,\omega}^{\text{FF}} |_{\Gamma_{\delta H}} W_{1,\omega}$ where $W_{1,\omega}$ is solution of:

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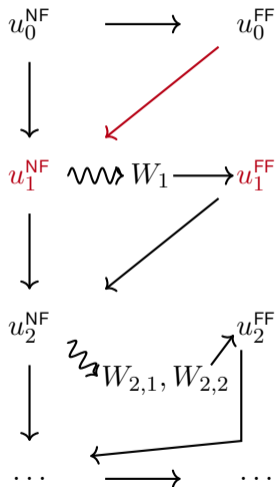
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Boundary Condition

$$u_{1,\omega}^{\text{NF}}(\cdot, Y) \xrightarrow{Y \rightarrow \infty} 0 \implies u_{1,\omega}^{\text{FF}}|_{\Gamma_{\delta H}} = \partial_\nu u_{0,\omega}^{\text{FF}}|_{\Gamma_{\delta H}} \lim_{Y \rightarrow \infty} W_{1,\omega}(\cdot, Y)$$

$$c_1 := \lim_{Y \rightarrow \infty} W_{1,\omega}(\cdot, Y) = \mathbb{E}[W_{1,\omega}(\cdot, H_\infty)]$$

Formal recursive construction

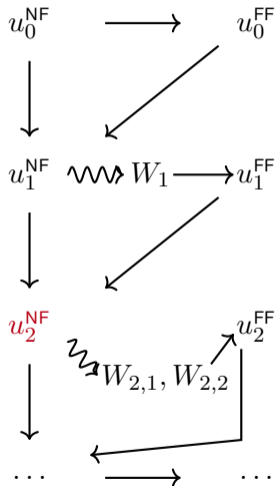


- u_1^{FF} satisfies:

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- $u_{1,\omega}^{\text{NF}}$ is written as $u_{1,\omega}^{\text{NF}} = (W_{1,\omega} - c_1 \chi_{Y>H}) \partial_\nu u_0^{\text{FF}}|_{\Gamma_{\delta H}}$

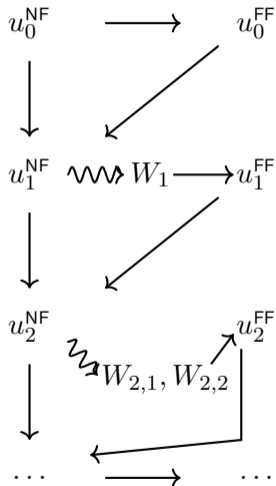
Formal recursive construction



■ Problem satisfied by $u_{2,\omega}^{\text{NF}}$:

$$\begin{cases} -\operatorname{div}_{X,Y}\left(\frac{1}{\mu_\omega}\nabla_{X,Y}u_{2,\omega}^{\text{NF}}\right) = \mathcal{F}_s\left(u_{1,\omega}^{\text{NF}}\right) & \text{in } \mathcal{B} \\ [u_{2,\omega}^{\text{NF}}]_{\Gamma_H} = -u_{2,\omega}^{\text{FF}}|_{\Gamma_{\delta H}} & [\partial_Y u_{2,\omega}^{\text{NF}}]_{\Gamma_H} = -\partial_\nu u_{1,\omega}^{\text{FF}}|_{\Gamma_{\delta H}} \\ u_{2,\omega}^{\text{NF}} = 0 & \text{on } \Sigma \end{cases}$$

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Du to the lack of higher decrease of the solution, we generally cannot define $u_{2,\omega}^{\text{NF}}$ and $u_{2,\omega}^{\text{FF}}$.

Effective models and error estimates

- In practice, we introduce truncated asymptotic expansions, which we approximate by functions satisfying effective impedance conditions at different orders

$$u_{\delta,N}^{\text{FF}} = \sum_{k=0}^N \delta^k u_k^{\text{FF}}$$

$$\text{Order 1 } u_{\delta,0}^{\text{FF}} = 0 \text{ on } \Gamma_{\delta H} \longrightarrow u_{\delta,0} = u_0^{\text{FF}}$$

$$\text{Order 2 } \partial_\nu u_{\delta,1}^{\text{FF}} + \frac{1}{\delta c_1} u_{\delta,1}^{\text{FF}} = O(\delta) \text{ on } \Gamma_{\delta H} \longrightarrow \partial_\nu u_{\delta,1} + \frac{1}{\delta c_1} u_{\delta,1} = 0 \text{ on } \Gamma_{\delta H}$$

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Effective models and error estimates

Error Estimates

For H big enough and every compact set K in $\mathbb{R}^2 \setminus \mathcal{D}_0$ such that $\text{dist}(K, \mathcal{D}_0) > 0$,

$$\sqrt{\mathbb{E} \left[\left\| u_{\omega, \delta} - u_{\delta, 0}^{\text{FF}} \right\|_{H^1(K)}^2 \right]} = O(\delta)$$

Under certain hypotheses on the randomness (Boucart, Fliss, and Giovangigli (2026)):

$$\sqrt{\mathbb{E} \left[\left\| u_{\omega, \delta} - u_{\delta, 1}^{\text{FF}} \right\|_{H^1(K)}^2 \right]} = O\left(\delta^{\frac{3}{2}} \sqrt{|\log(\delta)|}\right)$$

Computation of c_1 : random case

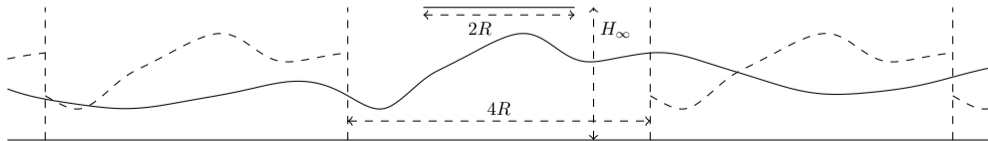
- $W_{1,\omega}$ is not computable in practice but using ergodicity we have

$$c_1 = \mathbb{E} [W_{1,\omega}(\cdot, H_\infty)] = \lim_{R \rightarrow \infty} \int_{-R}^R W_{1,\omega}(\cdot, H_\infty)$$

- Truncation of the computational domain of width $4R$ in the X direction by periodic conditions $\rightarrow W_{1,\omega,R}^\#$. Let $c_{1,R}^\# := \int_{-R}^R W_{1,\omega,R}^\#(\cdot, H_\infty)$

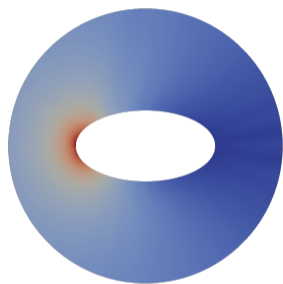
- We can estimate the error made by considering $W_{1,\omega,R}^\#$ instead of $W_{1,\omega}$ (relying on the work of A. Gloria)

$$\sqrt{\mathbb{E} [|c_1 - c_{1,R}^\#|^2]} \leq \underbrace{\sqrt{\mathbb{E} \left[\left| c_1 - \int_{-R}^R W_{1,\omega}(\cdot, H_\infty) \right|^2 \right]}}_{\lesssim \frac{1}{\sqrt{R}} \text{ under quantitative assumptions on the ergodicity}} + \underbrace{\sqrt{\mathbb{E} \left[\left| \int_{-R}^R W_{1,\omega}(\cdot, H_\infty) - c_{1,R}^\# \right|^2 \right]}}_{\lesssim \frac{\log(R)^2}{\sqrt{R}} \text{ under regularity hypothesis on } \mu_\omega}$$

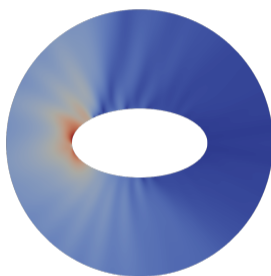
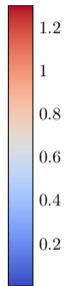


Error plots

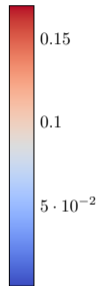
$$\kappa_0 = 10, \mu^{-1} = 1/3, \varepsilon = 1.4, \delta = 0.026$$



$$|u_\delta - u_{\delta,0}|$$



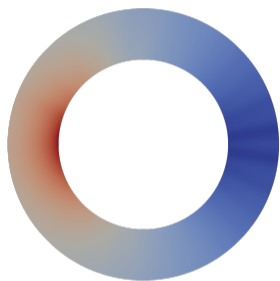
$$|u_{\omega,\delta} - u_{\delta,1}|$$



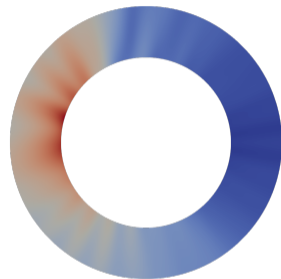
→ Implementation with Xlife++

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$$|u_\delta - u_{\delta,0}|$$



$$|u_{\omega,\delta} - u_{\delta,1}|$$

→ Implementation with Xlife++

Convergence plot

- Monte-Carlo method for calculating the expectation
- Reference solutions must be computed for a large number of realizations

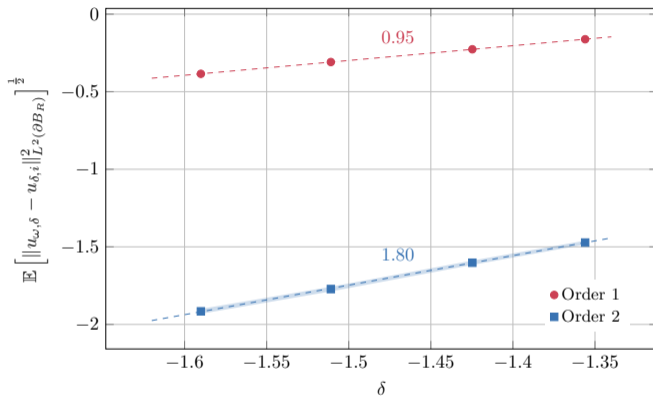
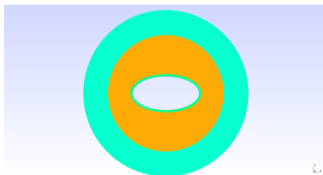


Figure: $L^2(\Omega, L^2(\partial B_R))$ -norm of the error between u_δ and $u_{\delta, i}$ for $R \gg 1$ ($N_{simu} = 220$)

Conesphere

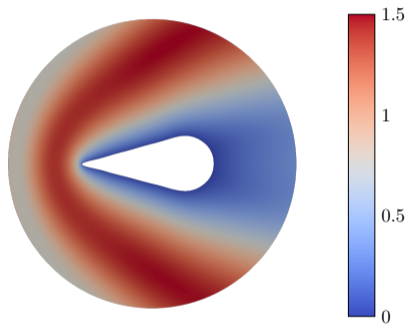


Figure: $|u_\delta|$, $\kappa_0 = 4$, $\mu^{-1} = 1/3$, $\varepsilon = 1.4$, $\delta = 0.026$

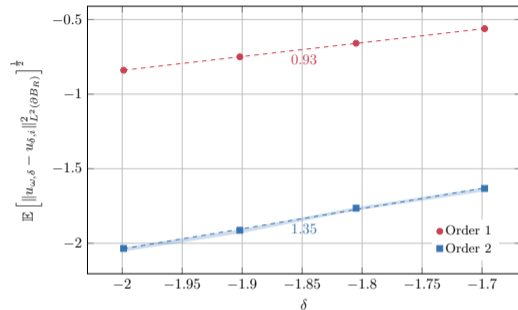
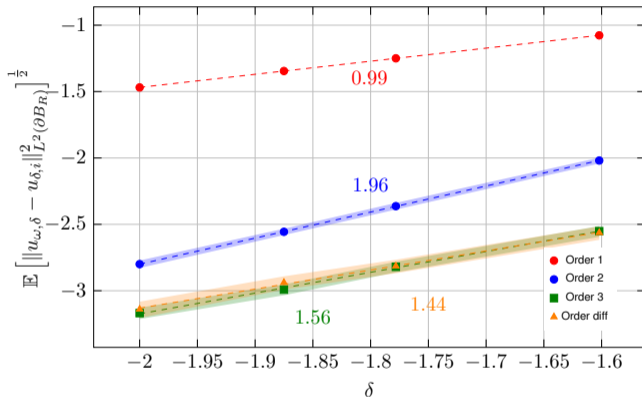
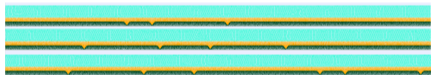


Figure: $L^2(\Omega, L^2(\partial B_R))$ -norm of the error between u_δ and $u_{\delta,i}$ for $R \gg 1$ ($N_{simu} = 130$)

Convergence rates have deteriorated due to the singular point (Vial (2003), Baudet (2024))

Higher orders of approximation

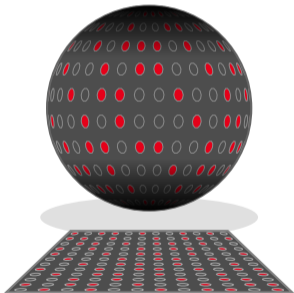
Simulations seem to show no higher order is defined, even with "kind" stochastic laws



$$\sqrt{\frac{1}{N} \sum_{i=1}^N \|u_{\omega_i, \delta} - u_{\omega_{N+i}, \delta}\|^2} \leq \sqrt{\frac{1}{N} \sum_{i=1}^N \|u_{\delta, 2}^{\text{FF}} - u_{\omega_i, \delta}\|^2} + \sqrt{\frac{1}{N} \sum_{i=N+1}^{2N} \|u_{\delta, 2}^{\text{FF}} - u_{\omega_i, \delta}\|^2}$$

Transition to 3D

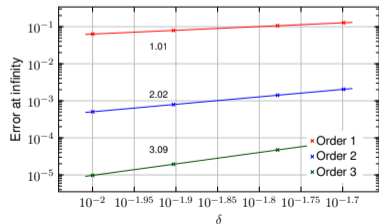
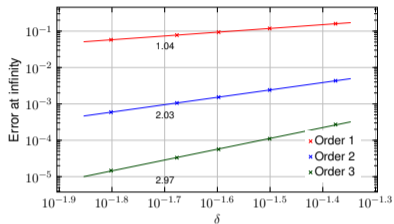
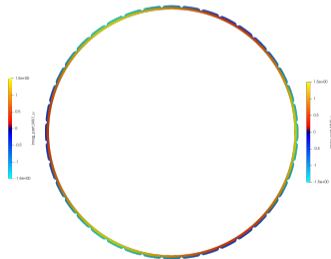
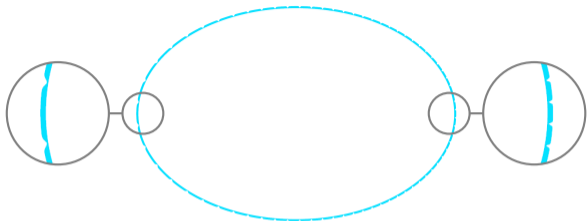
- Difficulty in defining stationarity on a 3D object (Chamaillard (2016))



- Introduction of deformations and patches
- Correctors' dependence on tangential variables
→ Calculation of correctors at each point on the surface
- High numerical cost

For a given realization, cell problems correspond to parametric PDEs. We use **basis reduction** techniques to reduce computation time (S. Hesthaven, Rozza, and Stamm (2016))

3D difficulties in 2D with basis reduction, $\mu \left(s; \frac{f(s)}{\delta}, \frac{\nu}{\delta} \right)$



$$\mu_{\delta,\omega}^{-1}\left(s; \frac{s}{\delta}, \frac{\nu}{\delta}\right)$$

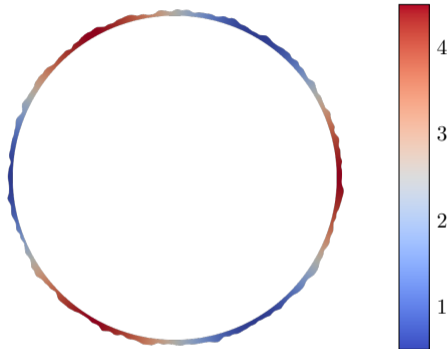


Figure: $\mu_{\delta,\omega}^{-1}$, $\kappa_0 = 1$, $\varepsilon = 1.4$, $\delta = 0.033$

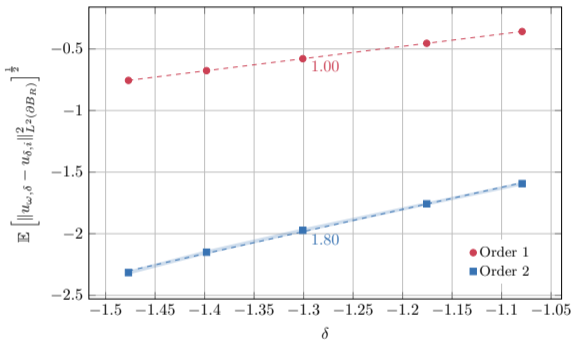


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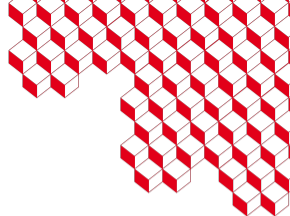
for $R \gg 1$ ($N_{simu} = 393$)

Conclusion and Perspectives



Problems under study

- ✓ 2D Helmholtz problem for any smooth geometry
- ✗ 3D Helmholtz problem for any smooth geometry
- ✗ 3D Maxwell problem for any smooth geometry
- ✗ Basis reduction for cell problems in high dimensions (Oseledets (2011))



Thanks for your attention!